## **Beyond the Tensor to Scalar ratio** CMB-S4 collaboration meeting

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# Assuming a detection of *r* has been made. Then what?

### Synergies and challenges Outline

- How do we know it is from inflation?
  - Look at statistics?
  - Cross-correlations?
- Challenges
  - Foregrounds and secondaries
  - `Intrinsic' signal
- For discussion: what if we don't detect a signal w S4?

#### **Beyond** *r* **Targets of interest**

- Prediction for PS:  $P_t(k) \propto r (k/k_*)^{n_t}$ 
  - In SFSR  $n_t < 0$ . Consistency check
  - But, in SFSR  $|n_t| \ll 1$ . Hard to do with CMB (see e.g. Dodelson 2014)
  - model dependent (see Kinney 2021)  $-N_{eff}$ ?
- Meerburg, Freese 2019)
  - For SFSR, all these are slow-roll suppressed (see e.g. Maldacena & Pimentel 2012).
  - massless, see e.g. Bordin et al 2016);
  - 2017)

• Can try using multi-messenger approach (Tania Regimbau, Robert Caldwell's talks, see also e.g. Meerburg et al 2015); but scaling really

• Beyond the PS, look at higher order correlation functions, e.g. bispectrum (see e.g. Muresuke 2014, Meerburg et al 2016, Duivenvoorden,

• Even when adding additional degrees of freedom there is a bound (Higuchi bound, mass of spin-2 mediator particle is bound, cant be

• Specifically, correlating e.g. a tensor ( $\gamma$ ) with two scalars ( $\zeta$ ) should have zero squeezed NGs even when adding a field. Caveat when breaking isometries of dS (e.g. solid inflation, Endlich et al 2012, Bordin et al 2018) or higher order partially massive particles (Baumann et al

• So to leading order, perhaps we can use the squeezed limit of tensor NGs to determine if gravitational waves are coming from `inflation' at all

#### **Tensor NGs Forecasts**

science book)

Shape: $\langle \mathcal{RR}\gamma \rangle$	Current	CMB-S4 goal	Conservative	CV-limited
$\langle BTT \rangle, \langle BTE \rangle, \langle BEE \rangle$				
$f_{ m sky}$	69%	3%	3%	100%
$\sigma(\sqrt{r}  ilde{f}_{ m NL}^{ m local})$	28	0.79	1.2	0.052
$\sigma(\sqrt{r}  ilde{f}_{ m NL}^{ m equil})$		16	24	1.7
$\sigma(\sqrt{r}  ilde{f}_{ m NL}^{ m ortho})$	•••	4.4	7.4	0.41

those in general, unfortunately, are harder to constrain

#### Forecasts show that we can do really well on squeezed limits (see S4 DSR,

NGs are therefore typically generated away from squeezed limit (equilateral);

## **Cross-correlations (2) Squeezed NGs**

- Assuming there exist tensor NGs. In the squeezed limit, can we cross correlate between different data?
- Example 1: primary CMB x spectral distortions
- Example 2: primary CMB x direct GWs

# Ex. 2: primary CMB x direct GWs

- Anisotropies in the energy density of primordial GW can be generated from squeezed  $\langle \gamma_{k_s} \gamma_{k_s} \zeta_{k_l} \rangle$ ,  $\langle \gamma_{k_s} \gamma_{k_s} \gamma_{k_l} \rangle$  (Adshead et al. 2020, Malhotra et al. 2021, Dimastrogiovanni et al. to appear)
- Long-short mode correlations leads to modulati density ( $\langle \gamma_{k_s} \gamma_{k_s} \rangle$ ) arising from different regions
- $\langle CMB GW \rangle$  probes ultra-squeezed configuration interferometer or PTA scales
- Needs significant enhancement of squeezed NG blue tensor spectrum  $n_t > 0$

ions of GW energy	CMB: $\overrightarrow{k}_{s}^{CMB}$ $\overrightarrow{-k}_{cMB}^{CMB}$
ations, $k_s^{GW}$ at	$\overrightarrow{k}_{s}^{SD}$
$Gs f_{NL} \gg 1$ , and	$-\overrightarrow{k}_{s}^{SD}$
GW:	$\overrightarrow{k}_{s}^{GW}$
	$-\overrightarrow{k}_{s}^{GW}$



## **Ex. 2: probing** $\langle \gamma \gamma \zeta \rangle$ from $\langle T - GW \rangle$

 $f_{NL} = f_{NI}^{\gamma\gamma\zeta}$  $\langle T - GW \rangle$  affected by  $\langle \gamma_{k_s} \gamma_{k_s} \zeta_{k_l} \rangle$ 

 $\delta^{GW}(k_{s},\hat{n}) = \int_{l_{s} \ll l_{s}} \frac{d^{3}k_{l}}{(2\pi)^{3}} e^{-i(\eta_{0}-\eta_{in})\hat{n}\cdot\vec{k}_{l}} f_{NL}(k_{l},k_{s})\zeta_{k_{l}}$ 

- Observations limited by low angular resolution of GW detectors ( $\ell_{\rm max} \sim 15 - 30$ ), need a high sensitivity network e.g. futuristic BBO
- Cross-correlations may also help to detect primordial anisotropies in presence of foregrounds
- Can also correlate  $\langle E GW \rangle$  from  $\langle \gamma_{k_s} \gamma_{k_s} \zeta_{k_l} \rangle$ , and nonzero  $\langle B - GW \rangle$  from  $\langle \gamma_{k_s} \gamma_{k_s} \gamma_{k_l} \rangle$  could hint to parity violation...



...more in Dimastrogiovanni, Fasiello, Malhotra, Meerburg, Orlando, to appear





Challenges

### Challenges for a detection **General limitations**

- Sources that look the same/similar (bias)
  - Statistically
  - Spectrally
- Sources that add variance (noise)
- Typically, both occur
- To mitagate:
  - Identify and model / project / de-source
- \*note that that we treat these as noise but these are also signal

### The CMB bispectrum Example

- Sources that look the same/similar (bias)
  - E.g. ISW-lensing (see Hill 2018, Coulton et al in prep)
- Sources that add variance (noise)
  - E.g. lensing (but in principle all above sources as well, see Coulton et al 2019)
- Galactic foregrounds; however here we will likely rely on simulations to check if they • contain statistics that is similar to signal; obviously cleaning the data, as we do for the PS, will be critical
- Note that higher order statistics in principle have the advantage that there are more dof, which benefits our ability to distinguish it from signal

## The CMB bispectrum Foregrounds (temperature only)



-ICMB-ICMB		
-ICMB-CIB	 Local	
-ICMB-kSZ	 Orthogonal	
-ICMB-tSZ	 Equilateral	
-ICMB-ISW		
CIB-ICMB		
CIB-CIB		
JB-kSZ		
JIB-tSZ		
CIB-ISW		
(SZ-ICMB		
(SZ-CIB		
(SZ-kSZ		
		Credits: Will Coulton

### The CMB bispectrum Intrinsic bispectrum

- Besides primordial and secondary sources, the CMB will also contain intrinsic bispectra, simply due to non-linear evolution of perturbations
- These could also be possible sources of confusion (and extra variance);
- Good news is that while they could be detectable with upcoming surveys (see Coulton 2021), they likely would not interfere with search for primordial NGs

#### Detectability w S4 $10^{1}$ **Evolution and Scat.** Quadratic Parity Even $10^{0}$ SNR $10^{-1}$ $10^{-2}$ -10<sup>3</sup> $10^{1}$ 10<sup>2</sup>

Credits: Will Coulton



## **Discussion and conclusions**

- If we detect r we should
  - Confirm it is from inflation
  - Look for statistics beyond the PS
- correlations (general synergies)
- Think more how to practically constrain GWs using large scale structure
- If we don't detect r?
  - Could still constrain (Maldacena) consistency relation
  - 2016)
  - Obviously, constraining trispectra will open up a new can of worms, machine learning?



• Challenges are well characterized for CMB only measurement, but we should think more about those for cross

• Could also look for trispectra which could potentially probe spin-2 fields (and higher) (see e.g. Bordin et al

## **Cross-correlations (1) Curl lensing**

- For sake of confidence, can we confirm the primordial nature of the GWs using other tracers?
- In the large scale structure, very challenging to 'constrain' tensor modes. (See e.g. Masui & Pen 2012, Schmidt et al 2013, Chisari et al 2014, Biagetti & Orlando 2020, `fossil' effects are promising)
- One example is curl lensing; presence of large scale primordial GWs can induce lensing signal with odd parity structure.
- In principle detectable; would provide proof of existence of large scale gravitational waves (Sheere, van Engelen, Meerburg, Meyers 2016)



### Intermezzo Why is it hard to constrain non-local NG in the CMB?

- First, tensors decay
- Second, and this is general for scalar/tensor NG in the CMB, on small scales NG are severely affected by blurring (See Kalaja, Meerburg, Pimentel & Coulton 2020)
- As a result, the improvement on NGs of these types does not improve as mode-counting
- Interestingly, higher n-point functions can exceed mode counting (e.g. trispectrum)





# Ex. 1: primary CMB x spectral distortions

- Spectral distortions are generated by the injection of energy from the dissipation of acoustic waves in the photon-baryon fluid.
- They are quadratic in primordial perturbations:  $\mu, y \sim \zeta^2, \gamma^2$
- Probe scales smaller than primary CMB.
- $\langle CMB \mu \rangle$ : sensitivity to very squeezed NGs
- Previous work considered scalar NGs with  $\langle T\mu \rangle$ ,  $\langle E\mu \rangle$ . (see e.g. Pajer and Zaldarriaga 2012, Emami et. al. 2015, Shiraishi et. al. 2015, Ota 2016, Ravenni et al. 2017, Cabass et. al. 2018)
- In the cosmic variance limit,  $\sigma(f_{NL}^{\rm loc}) \ll 1$





## **Ex. 1: tensor NGs from** $\langle CMB - \mu \rangle$

- $\gamma vs \zeta$  :  $\mu$  transfer function:
  - pro:  $\gamma$  transfer function probe a larger window of scales than  $\zeta$
  - con:  $\gamma$  transfer function is 5 orders of magnitude smaller than  $\zeta$
- Net effect: detecting squeezed  $\langle \gamma_l \gamma_s \gamma_s \rangle$ ,  $\langle \zeta_l \gamma_s \gamma_s \rangle$  is going to be challenging, any signal is obscured by  $\langle \gamma_l \zeta_s \zeta_s \rangle$ ,  $\langle \zeta_l \zeta_s \zeta_s \rangle$
- A large independent amplification on  $\gamma$  is needed, no viable models currently in literature (Orlando, Meerburg, Patil, to appear)
- On squeezed  $\langle \gamma_l \zeta_s \zeta_s \rangle$ :
  - probed by  $\langle B\mu 
    angle$
  - Signal is vanishing if bispectrum is isotropic (similar to  $\langle BT \rangle$  and  $\langle BE \rangle$ )
  - Need to introduce primordial anisotropies
  - Off diagonal  $\langle B\mu \rangle$  would be sourced by anisotropic NGs



Orlando, Meerburg, Patil, to appear